

The Jacobian Conjecture, a reduction of the degree via a Combinatorial Physics approach

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arXiv:1411.6558[math.AG], *Annales Henri Poincaré* (2016)

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Jacobian Conjecture

strikingly simple and natural conjecture

(a metro/tram ticket (size) conjecture)

"high school algebra"

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"high school algebra"

O. Keller, *Monats. Math. Phys.* (1939)

(for $n = 2$ and polynomials with integral coefficients)

Jacobian Conjecture (JC_n):

Let $n \geq 1$. If a polynomial function $F : \mathbb{C}^n \rightarrow \mathbb{C}^n$ has Jacobian determinant which is a non-vanishing constant, then the function F has a polynomial inverse.

Example.

$n = 2$, $F(z_1, z_2) = (z_1 + z_2^3, z_2)$ and $F^{-1}(z_1, z_2) = (z_1 - z_2^3, z_2)$.

$$\text{Jacobian} = \det \begin{pmatrix} \frac{dF_1}{dz_1} & \frac{dF_1}{dz_2} \\ \frac{dF_2}{dz_1} & \frac{dF_2}{dz_2} \end{pmatrix} = \det \begin{pmatrix} 1 & 3z_2^2 \\ 0 & 1 \end{pmatrix} = 1$$

Relations to the Dixmier conjecture

J. Dixmier, *Bull. Soc. Math. France* (1968)

Dixmier conjecture (DC_n):

Any endomorphism of the n -th Weyl algebra (the algebra of polynomial differential operators in n variables) is invertible.

$$DC_n \Rightarrow JC_n$$

A. Belov-Kanel and M. Kontsevitch, *Moscow Math. J.* (2007)

The Jacobian conjecture in dimension $2n$ implies the Dixmier conjecture in rank n .

Back to the Jacobian conjecture - faulty proofs

- before 1982 (see H. Bass *et. al.*, *Bull. Amer. Math. Soc.*)
 - W. Engel (*Math. Ann.* ('55)) claimed to prove the case $n = 2$.
A. Vitushkin (1975) published 2 essential errors
 - B. Segre published 3 incomplete proofs ('56,'57,'60) ; Canals and Lluís ('70) noted an error. Abhyankar and Moh pointed out a fault in Segre's proof and also in Canals and Lluís's correction
 - Gröbner proposed a proof in '61. Zariski pointed out that the argument is faulty
 - Oda in '80 proposed a proof - false
 - *etc.*
- after 1982 : other faulty proofs ...

A. van den Essen, "Polynomial automorphisms and the Jacobian conjecture", Birkhäuser (2000)

Internet blog: "How not to prove the Jacobian conjecture" :)

Back to the Jacobian conjecture - some notations

$$J_F(z) = \left(\frac{d}{dz_i} F_j(z) \right)_{1 \leq i, j \leq n}$$

\mathcal{P}_n - the set of polynomial systems F (all its coordinate functions F_j ($j = 1, \dots, n$) are polynomials)

$$\begin{aligned}\mathcal{J}_n^{\text{lin}} &:= \{F \in \mathcal{P}_n \mid \det J_F(z) = c \in \mathbb{C}^\times\}, \\ \mathcal{J}_n &:= \{F \in \mathcal{P}_n \mid F \text{ is invertible}\}.\end{aligned}$$

Jacobian Conjecture (JC_n):

$$\mathcal{J}_n^{\text{lin}} = \mathcal{J}_n \quad \forall n.$$

Some more notations

$$\deg(F) := \max_j \deg(F_j(z)),$$

$$\mathcal{P}_{n,d} := \{F \in \mathcal{P}_n \mid \deg(F) \leq d\}$$

$$\mathcal{J}_{n,d}^{\text{lin}} := \{F \in \mathcal{P}_n \mid \det J_F(z) = c \in \mathbb{C}^\times, \deg(F) \leq d\}$$

$$\mathcal{I}_{n,d} := \{F \in \mathcal{P}_n \mid F \text{ is invertible, } \deg(F) \leq d\}$$

Some (substantial?) progress

- 1 theorem for the quadratic case

Theorem

(S. Wang, *J. Alg.* (1980))

$$\mathcal{J}_{n,2}^{\text{lin}} = \mathcal{J}_{n,2} \quad \forall n.$$

- 2 reduction theorem to the cubic case

Theorem

(H. Bass et. al., *Bull. Am. Math. Soc.* (1982))

$$\mathcal{J}_{n,3}^{\text{lin}} = \mathcal{J}_{n,3} \quad \forall n \quad \implies \quad \mathcal{J}_n^{\text{lin}} = \mathcal{J}_n \quad \forall n.$$

$F_i(z) = z_i + \text{homogeneous pol. of degree 3.}$

A remark on H. Bass *et. al.* proof

The proof of H. Bass *et. al.* involves manipulations under which the dimension n is increased, thus this proof does *not* imply the corresponding statement without the “ $\forall n$ ” quantifier, i.e. that

$$\mathcal{J}_{n,3}^{\text{lin}} = \mathcal{J}_{n,3} \Rightarrow \mathcal{J}_n^{\text{lin}} = \mathcal{J}_n.$$

Our result - a further reduction of the degree; notations

For $n' \leq n$ and $F \in \mathcal{P}_{n,d}$, we write

$$z = (z_1, z_2)$$

and

$$F = (F_1, F_2)$$

to distinguish components in the two subspaces

$$\mathbb{C}^{n'} \times \mathbb{C}^{n-n'} \cong \mathbb{C}^n.$$

We set

$$R(z_2; z_1) = F_2(z_1, z_2),$$

emphasizing that, in R , we consider z_2 as the variables in a polynomial system, and z_1 as parameters.

The invertibility of R , denoted by $R(\cdot; z_1) \in \mathcal{J}_{n-n',d}$, for a fixed z_1 , means that there exists a pol. R^{-1} with variables $y_2 \in \mathbb{C}^{n-n'}$, and depending on z_1 , s. t.

$$\forall z_2 \in \mathbb{C}^{n-n'}, \quad R^{-1}(R(z_2; z_1); z_1) = z_2.$$

A few more notations

We define the subspaces of $\mathcal{P}_{n,d}$:

$$\mathcal{J}_{n,d;n'} := \{F \in \mathcal{P}_{n,d} \mid R(\cdot; z_1) \in \mathcal{J}_{n-n',d} \forall z_1 \in \mathbb{C}^{n'} \\ \text{and } F^{-1} \text{ restricted to } \mathbb{C}^{n'} \times \{0\} \text{ is in } \mathcal{P}_{n'}\}$$

$$\mathcal{J}_{n,d;n'}^{\text{lin}} := \{F \in \mathcal{P}_{n,d} \mid R(\cdot; z_1) \in \mathcal{J}_{n-n',d} \forall z_1 \in \mathbb{C}^{n'} \\ \text{and } (\det J_F)(z_1, R^{-1}(0, z_1)) = c \in \mathbb{C}^\times, \forall z_1 \in \mathbb{C}^{n'}\}$$

generalizations of $\mathcal{J}_{n,d}$ and resp. $\mathcal{J}_{n,d}^{\text{lin}}$

linear subspace of dimension $n - n'$ (the last $n - n'$ variables) on which z vanishes

QFT-inspired choices - they should become clear in the sequel

Our result - a further reduction of the degree

A. de Goursac et. al., *Annales Henri Poincaré* (2016)

Theorem

For $n \in \mathbb{N}$ and $d \geq 3$, there exists an injective map

$\Phi : \mathcal{P}_{n,d} \rightarrow \mathcal{P}_{n(n+1),d-1}$ satisfying

$$\Phi(\mathcal{J}_{n,d}^{\text{lin}}) \equiv \mathcal{J}_{n(n+1),d-1;n}^{\text{lin}} \cap \text{Im}(\Phi); \quad \Phi(\mathcal{J}_{n,d}) \equiv \mathcal{J}_{n(n+1),d-1;n} \cap \text{Im}(\Phi),$$

where $\text{Im}(\Phi) = \Phi(\mathcal{P}_{n,d})$.

A first consequence

Combining Bass *et. al.* theorem and the theorem above, the full Jacobian Conjecture reduces to the question whether

$$\mathcal{J}_{n(n+1),2;n}^{\text{lin}} \cap \text{Im}(\Phi) = \mathcal{J}_{n(n+1),2;n} \cap \text{Im}(\Phi).$$

- this question seems as difficult as the original Jacobian conjecture ...
- it involves **only a quadratic degree**, and this might simplify the resolution, in the light of Wang Theorem

A stronger version of JC

Conjecture

For all $n \geq n' \geq 0$, and all $d \geq 1$,

$$\mathcal{J}_{n,d;n'}^{\text{lin}} = \mathcal{J}_{n,d;n'}.$$

JC follows from the above conjecture

(0–dim.) Quantum Field Theory in a nutshell

A theory defined by means of a (functional) integral representation of the **partition function**, in which the **fields** are linearly coupled to **sources**;

from this, all the **correlation functions** of the respective physical system can be obtained by (functional) differentiation

A. Abdesselam, *Sém. Loth. Comb.* (2002),

A. Tanasă, *Sém. Loth. Comb.* (2012)

0–dim. Quantum Field Theory in a nutshell

Usually in QFT, the **fields** φ_i are functions of space(-time) (\mathbb{R}^D)

$D = 0$

($D \neq n$, the dimension of the linear system $F(z)$)

the scalar field φ_i is not a function of space-time

(there is no space-time)!

φ_i is a (real or complex) variable

partition function (generating function)

$$Z = \int_{\mathbb{R}} d\varphi e^{-\frac{1}{2}\varphi^2 + \frac{\lambda}{4!}\varphi^4}.$$

λ - the coupling constant

the quadratic part + interaction non-quadratic (here quartic) part

In 0–dim. QFT, the functional integral become usual (real or complex) integrals!

Combinatorial QFT

One (still) needs to evaluate integrals of type

$$\frac{\lambda^n}{n} \int d\varphi e^{-\varphi^2/2} \left(\frac{\varphi^4}{4!} \right)^n .$$

one can (still) use standard QFT techniques:

$$Z_0(J) := \int d\varphi e^{-\varphi^2/2 + J\varphi}$$

J - the source

computations of $(2k)$ -point correlation functions:

$$\int d\phi e^{-\phi^2/2} \phi^{2k} = \frac{\partial^{2k}}{\partial J^{2k}} \int d\varphi e^{-\varphi^2/2 + J\varphi} \Big|_{J=0} = \frac{\partial^{2k}}{\partial J^{2k}} e^{J^2/2} \Big|_{J=0} .$$

QFT - perturbation theory and Feynman graphs

perturbation theory - formal series in λ

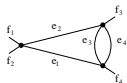
→ (abstract) Feynman graphs and Feynman integrals
(use of Wick Theorem)

A. Zvonkine, *Math. and Computer Modelling* (1997)

the quadratic part → edges

the interaction part → vertices

example:



(related to the physical information of a theory - interactions of elementary particle (in colliders a. s. o.))

0-dimensional QFT - interesting "laboratories" for testing theoretical physics tools

V. Rivasseau and Z. Wang, *J. Math. Phys.* (2010)

The intermediate field method in QFT - the idea

idea: introducing a new field, σ , to rewrite the interaction

the degree of the interaction has been reduced!

example: φ^6 model

$$Z(\lambda) = \int \frac{d\varphi}{\sqrt{2\pi}} e^{-\frac{1}{2}\varphi^2} e^{-\lambda\varphi^6} = \int \frac{d\varphi}{\sqrt{2\pi}} e^{-\frac{1}{2}\varphi^2} \int \frac{d\sigma}{\sqrt{2\pi}} e^{-\frac{1}{2}\sigma^2} e^{i\sqrt{2\lambda}\varphi^3\sigma}.$$

JC as a QFT model - the Abdesselam-Rivasseau model

A. Abdesselam, *Annales H. Poincaré* (2013)

$F \in \mathcal{P}_{n,d}$.

$$F_i(z) = z_i - \sum_{k=2}^d \sum_{j_1, \dots, j_k=1}^n w_{i,j_1 \dots j_k}^{(k)} z_{j_1} \dots z_{j_k} =: z_i - \sum_{k=2}^d W_i^{(k)}(z),$$

for $i \leq n$ and $w_{i,j_1 \dots j_k}^{(k)}$ some coefficients (the coupling constants)

JC as a QFT model - the Abdesselam-Rivasseau model

A. Abdesselam, *Annales Henri Poincaré* (2013)

the partition function

$$Z(J, K) = \int_{\mathbb{C}^n} d\varphi d\varphi^\dagger e^{-\varphi^\dagger \varphi + \varphi^\dagger \sum_{k=2}^d W^{(k)}(\varphi) + J^\dagger \varphi + \varphi^\dagger K},$$

where J, K are vectors in \mathbb{C}^n (the sources)

measure: $d\varphi d\varphi^\dagger := \prod_{i=1}^n \frac{d\operatorname{Re}\varphi_i d\operatorname{Im}\varphi_i}{\pi}$

$\varphi^\dagger K := \sum_{i=1}^n \varphi_i^\dagger K_i$, a. s. o.

setting the coupling constants to zero (free theory), the partition function is calculated by Gaussian integration:

$$\int_{\mathbb{C}^n} d\varphi d\varphi^\dagger e^{-\varphi^\dagger \varphi + J^\dagger \varphi + \varphi^\dagger K} = e^{J^\dagger K}.$$

very particular combinatorics of this QFT model

JC as a QFT model - the Abdesselam-Rivasseau model

- The partition function Z coincides with the inverse of the Jacobian
- The inverse G of F corresponds to the (standard) 1-point correlation function:

$$G_i(u) = \frac{\int_{\mathbb{C}^n} d\varphi d\varphi^\dagger \varphi_i e^{-\varphi^\dagger \varphi + \varphi^\dagger \sum_{k=2}^d W^{(k)}(\varphi) + \varphi^\dagger u}}{\int_{\mathbb{C}^n} d\varphi d\varphi^\dagger e^{-\varphi^\dagger \varphi + \varphi^\dagger \sum_{k=2}^d W^{(k)}(\varphi) + \varphi^\dagger u}} \quad (1)$$

The sets of polynomial functions involved in JC can be rephrased in this framework:

$$\mathcal{J}_{n,d}^{\text{lin}} = \{F \in \mathcal{P}_{n,d} \mid Z(0, u) = 1, \forall u \in \mathbb{C}^n\},$$

$$\mathcal{J}_{n,d} = \{F \in \mathcal{P}_{n,d} \mid G_i(u) \text{ given by (1) is in } \mathcal{P}_n\}.$$

The intermediate field method for the JC QFT model

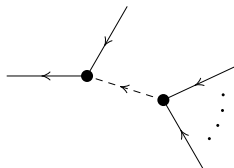
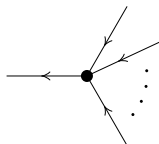
- This will reduce the degree d of F .
- We will add n^2 "intermediate fields" σ_{ij} to the model.

$i, j = 1, \dots, n$

- intermediate field identity

Using the general formula of Gaussian integration, one has:

$$e^{(\varphi_i^\dagger \varphi_j) \left(\sum_{j_2, \dots, j_d=1}^n w_{i,j,j_2 \dots j_d}^{(d)} \varphi_{j_2} \dots \varphi_{j_d} \right)}$$
$$= \int_{\mathbb{C}^{n^2}} d\sigma_{i,j} d\sigma_{i,j}^\dagger e^{-\sigma_{i,j}^\dagger \sigma_{i,j} + \sigma_{i,j}^\dagger \left(\sum_{j_2, \dots, j_d=1}^n w_{i,j,j_2 \dots j_d}^{(d)} \varphi_{j_2} \dots \varphi_{j_d} \right) + (\varphi_i^\dagger \varphi_j) \sigma_{i,j}}$$



The intermediate field method for the JC QFT model

use the intermediate field identity for each pair (i, j) , in the partition function $Z(J, K)$ of the JC QFT model with n dimensions and degree d , in order to re-express the monomials of degree d in the fields φ

$$\begin{aligned} \Rightarrow Z(J, K) = & \int_{\mathbb{C}^n} d\varphi d\varphi^\dagger \int_{\mathbb{C}^{n^2}} d\sigma d\sigma^\dagger e^{-\varphi^\dagger \varphi + \varphi^\dagger \sum_{k=2}^{d-1} W^{(k)}(\varphi) + J^\dagger \varphi + \varphi^\dagger K} \\ & e^{\sum_{i,j=1}^n \left(-\sigma_{i,j}^\dagger \sigma_{i,j} + \sigma_{i,j}^\dagger \sum_{j_2, \dots, j_d=1}^n w_{i,j,j_2 \dots j_d}^{(d)} \varphi_{j_2} \dots \varphi_{j_d} + \varphi_i^\dagger \varphi_j \sigma_{i,j} \right)}. \end{aligned}$$

Setting this proper

- 1 We define the new **fields** $\phi \in \mathbb{C}^{n+n^2}$ by
$$\phi = (\varphi_1, \dots, \varphi_n, \sigma_{1,1}, \dots, \sigma_{1,n}, \dots, \sigma_{n,1}, \dots, \sigma_{n,n}).$$
- 2 We define the **coupling constants** \tilde{w} as:
 - for $k = d - 1$, we set $\tilde{w}_{i,j,j_2 \dots j_d}^{(d-1)} := w_{i,j,j_2 \dots j_d}^{(d-1)}$ and
 $\tilde{w}_{i,n+j,j_2 \dots j_d}^{(d-1)} = w_{i,j,j_2 \dots j_d}^{(d)}$ with $i, j, j_2, \dots, j_n \leq n$
 - for $k \in \{3, \dots, d - 2\}$, we set $\tilde{w}_{i,j,j_2 \dots j_k}^{(k)} := w_{i,j,j_2 \dots j_k}^{(k)}$ with
 $i, j, j_2, \dots, j_n \leq n$
 - for $k = 2$, we set $\tilde{w}_{i,j,j_2}^{(2)} := w_{i,j,j_2}^{(2)}$ and $\tilde{w}_{i,j,i \cdot n+j}^{(2)} = 1$ with
 $i, j, j_2 \leq n$.

The remaining coefficients of \tilde{w} are set to 0.

- 3 The **sources** are defined to be $\tilde{J} := (J, 0)$ and $\tilde{K} := (K, 0)$,
(the number of extra vanishing coordinates is n^2).

The resulting QFT model

- 1 The partition function:

$$Z(J, K) = \int_{\mathbb{C}^{n+n^2}} d\phi d\phi^\dagger e^{-\phi^\dagger \phi + \phi^\dagger \sum_{k=2}^{d-1} \tilde{W}^{(k)}(\phi) + \tilde{J}^\dagger \phi + \phi^\dagger \tilde{K}}$$

- 2 The 1-point correlation functions:

$$G_i(u) = \frac{\int_{\mathbb{C}^{n+n^2}} d\phi d\phi^\dagger \phi_i e^{-\phi^\dagger \phi + \phi^\dagger \sum_{k=2}^{d-1} \tilde{W}^{(k)}(\phi) + \phi^\dagger \tilde{u}}}{\int_{\mathbb{C}^{n+n^2}} d\phi d\phi^\dagger e^{-\phi^\dagger \phi + \phi^\dagger \sum_{k=2}^{d-1} \tilde{W}^{(k)}(\phi) + \phi^\dagger \tilde{u}}},$$

for $i \in \{1, \dots, n\}$.

So, what have we showed?

The partition function (resp. the 1–point correlation function) of the JC QFT model with dimension n and degree d is equal to the partition function (resp. the n first coordinates of the 1–point correlation function) of the model with dimension $n(n+1)$ and degree $d-1$, up to a redefinition of

- 1 the fields
- 2 the coupling constant $w \mapsto \tilde{w}$
- 3 the sources.

Since the partition function corresponds to the inverse of the Jacobian (resp. the 1–point correlation function corresponds to the formal inverse), *this gives a QFT proof of the theorem.*

However ...

This is not a proof of the Jacobian Conjecture (unfortunately)!

Alternative proof

In

A. de Goursac *et. al.* *Annales H. Poincaré* (2016)

algebraic (no-QFT-like) proof of our reduction result

various purely combinatorial approaches to JC were given:

- proposition of Joyal's combinatorial species as a tool
D. Zeilberger (1987)
- reformulation of the JC using trees
D. Wright (1999)
- reformulation of the JC using rooted trees
D. Singer, *Electron. J. Comb.* (2011)
- *etc.*

Conclusion and perspectives

- We have proved, using QFT-inspired techniques, a reduction theorem to the quadratic case for JC, up to the addition of a new parameter n' (related to the introduction of additional intermediate fields σ)
- immediate perspective: adaptation of Wang's proof to our modified quadratic case
- reformulation of Wang's proof in a QFT language
- revisit the Dixmier Conjecture from the perspective of non-commutative QFT

Thank you for your attention!

Vă mulțumesc pentru atenție!